

$\bullet a \sin(kx) \cdot \cos(\sigma t) = \frac{2A_2 D \sigma e^{kd}}{g} \cdot \cosh kd \cdot \sin(kx) \cdot \cos(\sigma t)$

$$2A_2 D e^{kd} = \frac{ag}{\sigma \cosh kd}$$

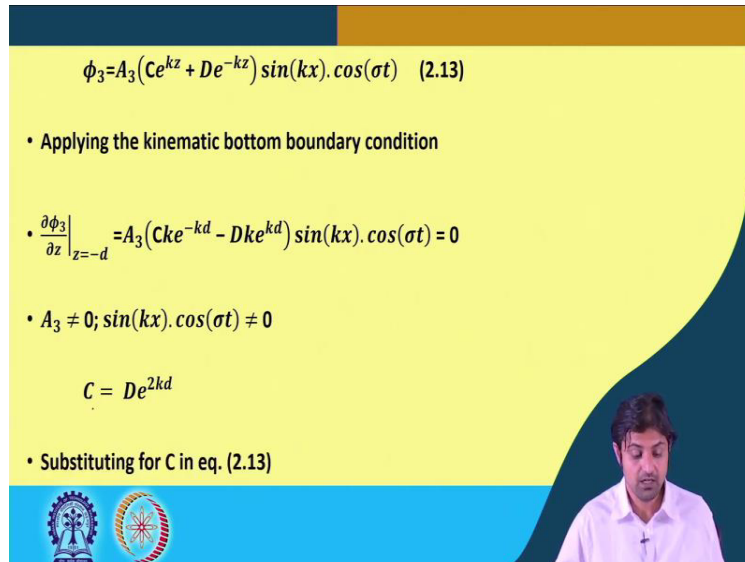
Substituting in eq. (2.11), we get

$$\phi_2 = \frac{ag}{\sigma} \frac{\cosh k(d+z)}{\cosh kd} \sin(kx) \cdot \sin(\sigma t) \quad (2.12)$$

Let us consider  $\phi_3$

I do not expect you to remember the derivation but the steps you must know the things like the specifying the governing equation, what are the governing equations governing equation is Laplace equation for the boundary conditions we utilize the Bernoulli's equation and the continuity equation for example, so, we have obtained phi 2 here.

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$$\phi_3 = A_3 (C e^{kz} + D e^{-kz}) \sin(kx) \cdot \cos(\sigma t) \quad (2.13)$$

$\bullet$  Applying the kinematic bottom boundary condition

$$\bullet \left. \frac{\partial \phi_3}{\partial z} \right|_{z=-d} = A_3 (C k e^{-kd} - D k e^{kd}) \sin(kx) \cdot \cos(\sigma t) = 0$$

$\bullet A_3 \neq 0; \sin(kx) \cdot \cos(\sigma t) \neq 0$

$$C = D e^{2kd}$$

$\bullet$  Substituting for C in eq. (2.13)

So, if we consider phi 3 and apply the same concepts that phi 3 = you know, we apply the dynamic boundary condition. This will give C equal to D into e to the power 2 k d.

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Substituting in eq. (2.14), we get

$$\phi_3 = \frac{-ag}{\sigma} \frac{\cosh k(d+z)}{\cosh kd} \sin(kx) \cdot \cos(\sigma t) \quad (2.15)$$


Let us consider  $\phi_4$

$$\phi_4 = A_4 (C e^{kz} + D e^{-kz}) \cos(kx) \cdot \sin(\sigma t) \quad (2.16)$$

Applying the kinematic bottom boundary condition

$$\left. \frac{\partial \phi_4}{\partial z} \right|_{z=-d} = A_4 (C k e^{-kd} - D k e^{kd}) \cos(kx) \cdot \sin(\sigma t) \longrightarrow C = D e^{2kd}$$

Substituting for C in eq. (2.16)



And same procedure is repeated for the dynamic free surface boundary condition and we get for  $\phi_3$  at term like this, you understand same procedure. So  $\phi_3$  we get  $-ag$  by  $\sigma \cosh kd + z$  similarly, we get  $\phi_3$ . And same procedure we do for  $\phi_4$ , for obtaining the values.

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$$\phi_4 = 2A_4 D e^{kd} \cosh k(d+z) \cos(kx) \cdot \sin(\sigma t) \quad (2.17)$$

And

$$\left. \frac{\partial \phi_4}{\partial t} \right|_{z=0} = 2A_4 D e^{kd} \cosh k(d+z) \cos(kx) \cdot \sin(\sigma t)$$

Assuming  $\eta = a \cos(kx) \cdot \cos(\sigma t)$  and applying eq (2.5)

We get  $2A_4 D e^{kd} = \frac{ag}{\sigma} \frac{1}{\cosh kd}$



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- Substituting in eq. (2.17), we get

$$\phi_4 = \frac{ag}{\sigma} \frac{\cosh k(d+z)}{\cosh kd} \cos(kx) \sin(\sigma t) \quad (2.18)$$

- Let us consider  $\phi_1$

$$\phi_1 = A_1 (C e^{kz} + D e^{-kz}) \cos(kx) \cos(\sigma t) \quad (2.19)$$

- Applying the kinematic bottom boundary condition

$$\left. \frac{\partial \phi_1}{\partial z} \right|_{z=-d} = 0 \quad C = D e^{2kd}$$


And what we get  $\phi_4$  is  $ag$  by  $\sigma$  a similar result we get and we repeat the same procedure again for  $\phi_1$ .

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- Substituting for  $C$  in eq. (2.19)


$$\phi_1 = 2A_1 D e^{kd} \cosh k(d+z) \cos(kx) \cos(\sigma t) \quad (2.20)$$

- assuming  $\eta = a \cos(kx) \sin(\sigma t)$  and applying equation (2.5)

We get

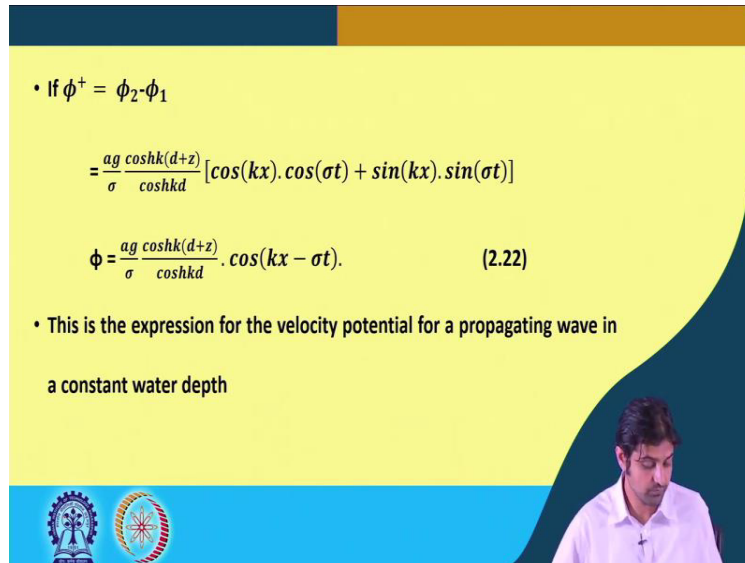
$$2A_1 D e^{kd} = \frac{-ag}{\sigma} \frac{1}{\cosh kd}$$

- Substituting in eq. (2.20), we get

$$\phi_1 = \frac{-ag}{\sigma} \frac{\cosh k(d+z)}{\cosh kd} \cos(kx) \cos(\sigma t) \quad (2.21)$$


Applying the kinematic bottom boundary condition first and then applying the dynamic. See I am skipping because it is essentially the same process.

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- If  $\phi^+ = \phi_2 - \phi_1$

$$= \frac{ag \cosh k(d+z)}{\sigma \cosh kd} [\cos(kx) \cdot \cos(\sigma t) + \sin(kx) \cdot \sin(\sigma t)]$$

$$\phi = \frac{ag \cosh k(d+z)}{\sigma \cosh kd} \cdot \cos(kx - \sigma t). \quad (2.22)$$

- This is the expression for the velocity potential for a propagating wave in a constant water depth

More important is this term so  $\phi_1$  we get is  $-ag$  by  $\sigma \cos h k d + z$  if you remember,  $\phi_2$  and  $\phi_4$  was positive with a positive sign  $\phi_4$  and  $\phi_1$  was with the negative sign after the derivation. Now you remember I said that the total velocity potential will be the summation of the two terms. So our velocity potential is going to be  $\phi_2 - \phi_1$ , or also in terms of  $\phi_3$  and  $\phi_4$  also so this becomes, so if you add to velocity potential, this was you know,  $\phi_2$ .

And, this one with a negative sign was  $\phi_1$  so, we get this so,  $\cos kx \cos \sigma t + \sin kx \sin \sigma t$  can be return as  $\cos kx - \sigma t$  using trigonometry. So, the final velocity potential with the value I mean the formula for which you are supposed to remember is this  $1 ag$  by  $\sigma \cos h k$  into  $d + z$  divided by  $\cos h k d$  into  $\cos kx - \sigma t$  this is the final velocity potential. So, now if you see we have determined  $\phi$  is the amplitude of the wave.

$\sigma$  is  $2\pi$  by  $t$   $k$  is  $2\pi$  by  $L$  everything you know. So, as a function of  $x$  we have found out the velocity potential here as I said this is the expression for the velocity potential for a propagating wave in a constant water depth. So, this is at one particular constant water depth, this is the velocity potential which we have derived from hydraulic assumptions and conditions.


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Since  $\eta = \frac{1}{g} \frac{\partial \phi}{\partial t} \Big|_{z=0}$

$$\eta = \frac{1}{g} \frac{ag \cosh k(d+z)}{\sigma \cosh kd} \cdot \sigma \sin(kx - \sigma t).$$

Hence  $\eta = a \sin(kx - \sigma t)$ . phase (2.23)

- ' $\eta$ ' is periodic in  $x$  and  $t$ . If we locate a point and traverse along the wave, such that, at all-time ' $t$ ' our position relative to the wave form remains fixed then the phase difference is zero or  $kx - \sigma t = \text{constant}$



Now, because we have got the velocity potential we can finally write the term for eta, we are going to check so, eta can be return as so,  $\frac{1}{g} \frac{\partial \phi}{\partial t}$  at  $z = 0$  will come out to be this one. So,  $ag$  can be return as so, if you just  $g$  and  $g$  will cancel  $\sigma$  and  $\sigma$  will cancel. So, at  $z = 0$ , this will also get canceled. So, simply we left with  $\eta = a \sin kx - \sigma t$ .

So, now this eta is periodic in  $x$  and  $t$ , if we locate a point and traverse along the wave says that at all time  $t$ , our position relative to the waveform remains fixed. So, basically this term  $kx - \sigma t$  from your earlier mechanics class you would know this is what this is the phase  $kx - \sigma t$ . So, with this phase we do an experiment what we do is if we locate a point and traverse along the waves such that at any time  $t$  our position relative to the waveform remains fixed.

Then this means that  $kx - \sigma t$  will be 0 phase difference is going to be 0 which means that  $kx - \sigma t$  is going to be constant. Because if phase will be constant phase difference is going to be 0.

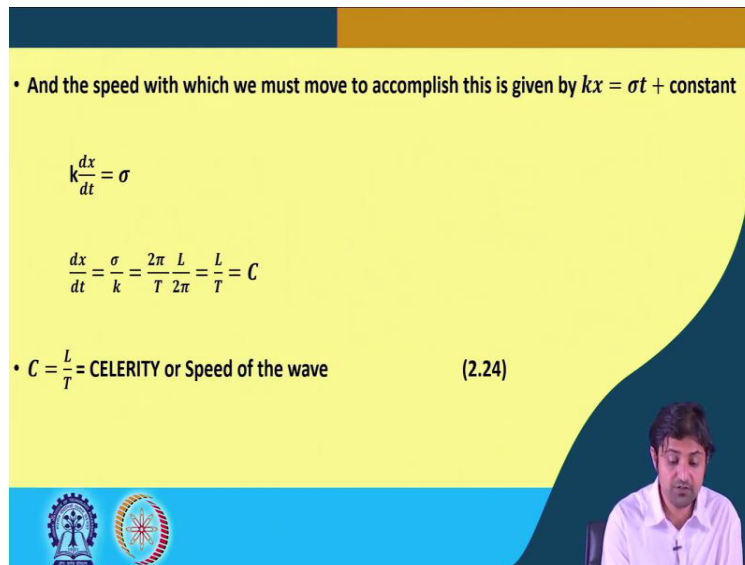
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- And the speed with which we must move to accomplish this is given by  $kx = \sigma t + \text{constant}$

$$k \frac{dx}{dt} = \sigma$$

$$\frac{dx}{dt} = \frac{\sigma}{k} = \frac{2\pi}{T} \frac{L}{2\pi} = \frac{L}{T} = C$$

- $C = \frac{L}{T} = \text{CELERITY or Speed of the wave}$  (2.24)



And the speed with which we must move to accomplish this will be given by  $kx - \sigma t = \text{constant}$  or  $kx = \sigma t + \text{constant}$ . On differentiating we can get this equation we differentiate we get  $k \frac{dx}{dt} = \sigma$  we differentiate with respect to time =  $\sigma$  or  $\frac{dx}{dt}$  can be written as  $\sigma$  by  $k$ , and  $\sigma$  was nothing but  $\frac{2\pi}{T}$  and  $k$  was nothing but  $\frac{2\pi}{L}$ . So, it becomes  $\frac{L}{T}$  and that is the  $C$  and  $C$  is the wave celerity.

So, to determine the velocity of the wave, this is the so, this is how the wave celerity that is length wave length by time period is the celerity and this is the basis of finding the way of celerity that if locate a point and traverse along the wave such that at all time  $t$  our position relative to the waveform remains fixed when this will happen when we are going to move with the speed of the wave. And that our speed and wave speed will be the same in that case and that is what we have found out that  $C = \frac{L}{T}$  is the celerity of the wave with which we must be moving.

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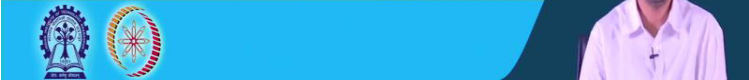
### Wave moving in negative 'x' direction

- If  $\phi^- = \phi_2 + \phi_1$ 

$$= \frac{-ag}{\sigma} \frac{\cosh k(kx + \sigma t)}{\cosh kd} \cdot \cos(kx + \sigma t).$$
- Since  $\eta = \frac{1}{g} \frac{\partial \phi}{\partial t} \Big|_{z=0}$ 

$$\eta = \frac{1}{g} \frac{-ag}{\sigma} \frac{\cosh k(d+z)}{\cosh kd} (-\sigma \sin(kx + \sigma t))$$

$$\eta = a \sin(kx + \sigma t)$$




Now, if there is a wave that is moving in x direction, then we can simply write the velocity potential  $\phi^- = \phi_2 + \phi_1$  not  $-\phi_1$  and then we will get nothing but with  $-\sin$  - ag by  $\sigma \cos h k x \sigma k x + \sigma t$  into  $\cos h k d$ . So, what you do is you do  $x = so$ , instead of  $k x + - \sigma t$  becomes  $\cos k x + \sigma t$ . So, this is a different waveform. So, in this case, since a  $\eta$  is

1 by g we will also get a similar you know, if there was it will remember it was  $k x - \sigma t$  it becomes  $k x + \sigma t$ . So, the phase is changing when the wave is traveling in the negative direction.

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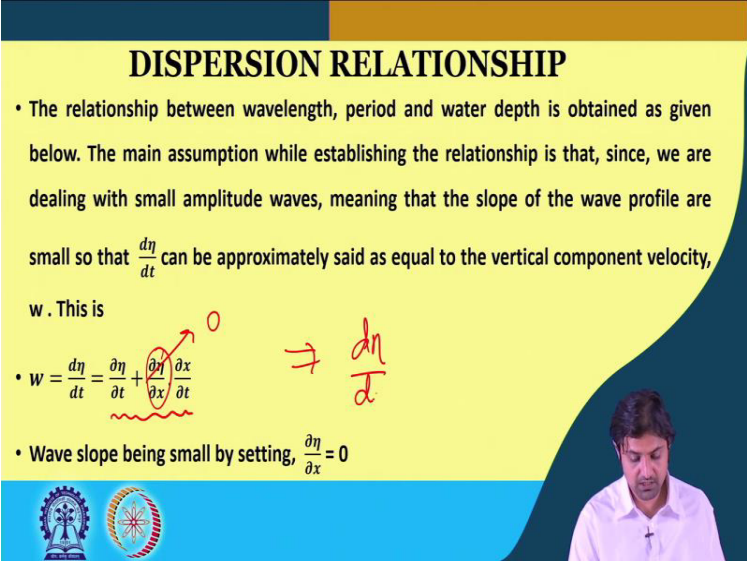
To obtain the celerity of the wave we have

$$kx + \sigma t = \text{constant}$$

$$\frac{dx}{dt} = \frac{-\sigma}{k} = \frac{-2\pi}{T} \frac{L}{2\pi} = \frac{-L}{T} = -C \quad (2.25)$$


Now, to obtain the celerity of the wave in this case we have  $kx + \sigma t = \text{constant}$  same procedure is repeated. So, it will become  $-c$  sonic wave negative means negative celerity in the negative  $x$  direction.

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**DISPERSION RELATIONSHIP**

- The relationship between wavelength, period and water depth is obtained as given below. The main assumption while establishing the relationship is that, since, we are dealing with small amplitude waves, meaning that the slope of the wave profile are small so that  $\frac{d\eta}{dt}$  can be approximately said as equal to the vertical component velocity,  $w$ . This is
- $w = \frac{d\eta}{dt} = \frac{\partial \eta}{\partial t} + \frac{\partial \eta}{\partial x} \frac{\partial x}{\partial t}$   $\Rightarrow \frac{d\eta}{dt}$
- Wave slope being small by setting,  $\frac{\partial \eta}{\partial x} = 0$

Now, one important thing after this derivation of the velocity potential is something called a dispersion relationship that is one of the core concept of the wave mechanics. So, the relationship between wavelength with period and water depth is obtained as given below for the dispersion relationship, the main assumption while establishing the relationship is that since we are dealing with small amplitude waves, meaning that the slope of wave profile are so, small that  $\frac{d\eta}{dt}$  can be approximately said equal to the vertical component of the velocity  $w$ .

So,  $w$  can be written as  $\frac{d\eta}{dt}$  or in differential form  $\frac{d\eta}{dt} = \frac{\partial \eta}{\partial t} + w$  is what  $\frac{d\eta}{dt}$  because that is the vertical velocity correct  $\frac{d\eta}{dt}$ . So, that if you apply the total you know, differentiation by you know, total differentiation, it will be  $\frac{d\eta}{dt} = \frac{\partial \eta}{\partial t} + \frac{\partial \eta}{\partial x} \frac{dx}{dt}$ . So, wave slope being small means  $\frac{\partial \eta}{\partial x} = 0$  which implies  $\frac{d\eta}{dt} = \frac{\partial \eta}{\partial t}$ .


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$w = \frac{\partial \eta}{\partial t}$  but  $w = -\frac{\partial \phi}{\partial z}$   
 Hence  $\frac{\partial \eta}{\partial t} = -\frac{\partial \phi}{\partial z}$  (2.26)

Differentiating the expression of  $\eta$  we get  $\frac{\partial \eta}{\partial t} = \frac{1}{g} \frac{\partial^2 \phi}{\partial t^2} \Big|_{z=0}$

Hence  $\frac{\partial \eta}{\partial t} = \frac{-A\sigma^2}{g} \cosh kd \cdot \cos(kx - \sigma t)$ . (2.27)



$W = \text{del } \eta \text{ del } t$  but  $w$  also equal to  $-\text{del } \phi \text{ del } z$  in form of velocity potential hence  $\text{del } \eta \text{ del } t$  equal to  $-\text{del } \phi \text{ del } z$  we know the velocity potential we know the equation of a  $\eta$ , if we do the differentiation of a  $\eta$  we get  $\text{del } t = \text{one}$ . So, if you do the differentiating the expression of  $\eta$ , what do we get?  $\text{Del } \eta \text{ by } \text{del } t = 1$  by  $g$  you remember, that term headset  $= 0$   $1$  by  $g$   $\text{del } \text{square } \phi \text{ by } \text{del } t \text{ square}$ .

Eta was  $1$  by  $g$   $\text{del } \phi \text{ del } t$  at  $z = 0$ . So, if you differentiate this you get  $\text{del } \eta \text{ del } t = \text{one}$  by  $g$   $\text{del } t \text{ square } \phi \text{ by } \text{del } t \text{ square}$  that  $= 0$  this is what it is written here so, we substitute  $\text{del } \eta \text{ del } t$  we get from here.


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*previous page*  
 • Where  $A = \frac{H g}{2 \sigma \cosh kd}$

$w = \frac{-\partial \phi}{\partial z} = -Ak \sinh kd \cdot \cos(kx - \sigma t)$ . (2.28)

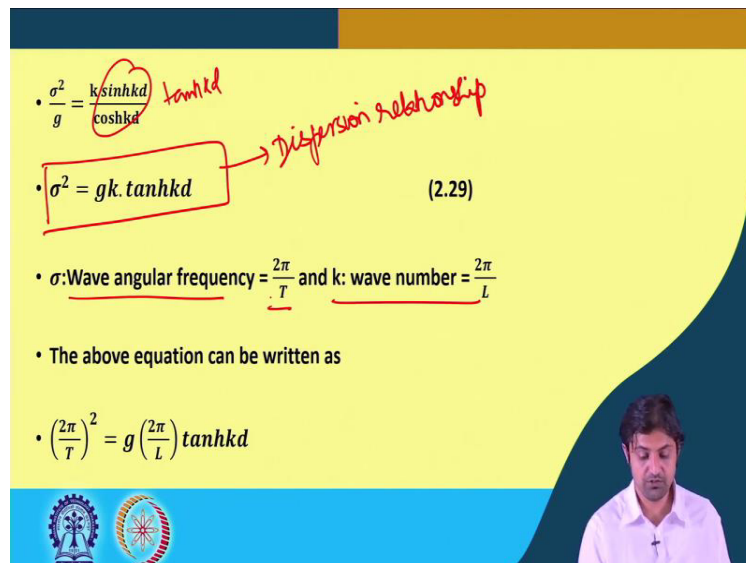
• Using the relation of Eq. (2.26), equating Eq. (2.27) to Eq (2.28), we get

$\frac{A\sigma^2}{g} \cosh kd \cdot \cos(kx - \sigma t) = Ak \sinh kd \cdot \cos(kx - \sigma t)$ .



Where  $\sigma$  = we write capital A is where A is here is  $H$  by 2 into  $g$  by  $\sigma$  divided one by  $\cos h k d$  and  $w$  is also  $-\frac{\partial \pi}{\partial z}$  it that is  $-Ak \sin h k d$  into  $\cos k x$  into  $\sigma t$ . And we equate this both these term  $A^2 \sigma^2$  from the previous equation with this one here from the previous page because both are same.

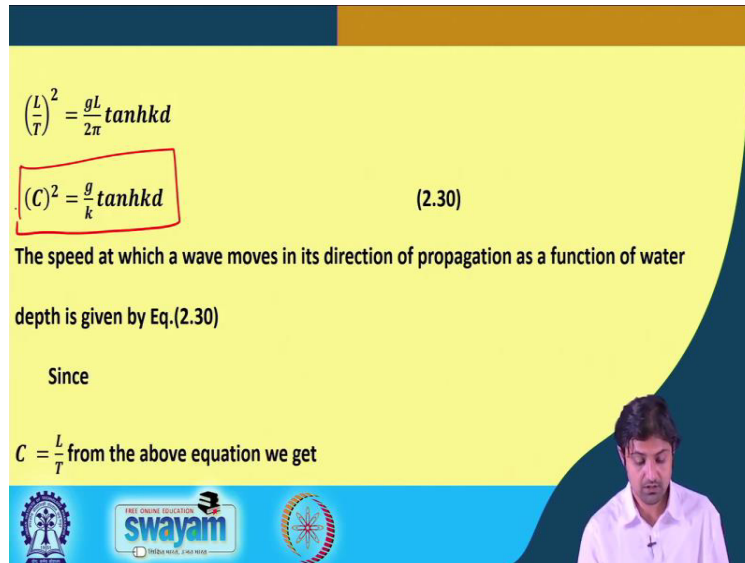
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$\sigma^2 = \frac{k \sinh kd}{\cosh kd} \tanh kd$   
 $\sigma^2 = gk \tanh kd$  (2.29) *→ Dispersion relationship*  
 $\sigma$ : Wave angular frequency  $= \frac{2\pi}{T}$  and  $k$ : wave number  $= \frac{2\pi}{L}$   
 • The above equation can be written as  
 $\left(\frac{2\pi}{T}\right)^2 = g \left(\frac{2\pi}{L}\right) \tanh kd$

And on canceling the common terms what we get is  $\sigma^2$  by  $g = k \sinh kd$  divided by  $\cosh kd$  and this is nothing but  $\tanh kd$ . So,  $\sigma^2$  can be written as  $gk \tanh kd$ . Here  $\sigma$  is the wave angular frequency equal to  $2\pi$  by  $T$  and  $k$  = wave number which  $= 2\pi$  by  $L$  as we have seen before and this above equation can be written as  $2\pi$  by  $T$  whole square  $= g$  into  $2\pi$  by  $L$  into  $\tanh kd$ . So, actually this is the famous dispersion relationship.

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$$\left(\frac{L}{T}\right)^2 = \frac{gL}{2\pi} \tanh kd$$

$$(C)^2 = \frac{g}{k} \tanh kd \quad (2.30)$$

The speed at which a wave moves in its direction of propagation as a function of water depth is given by Eq.(2.30)

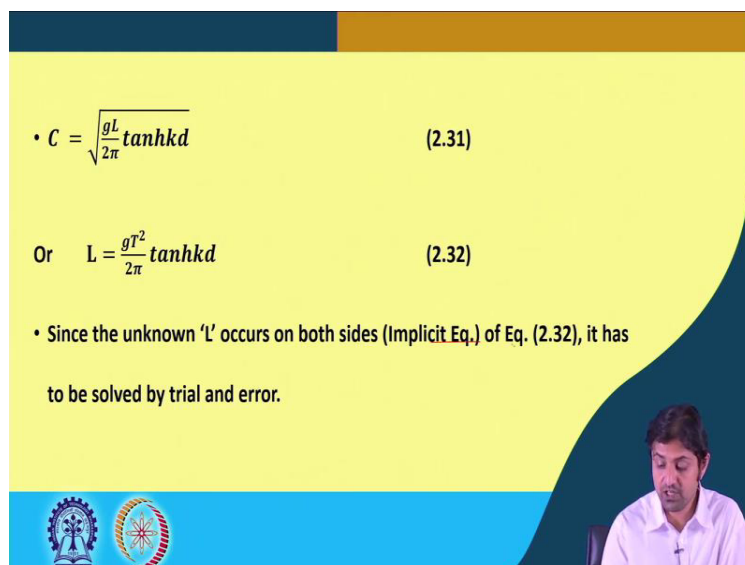
Since

$C = \frac{L}{T}$  from the above equation we get

swayam

So, but is substituting since sigma and k in terms we get L by T whole square = gL by 2 pi into tan h k d or C squared =, because it can be written as C square = g by k into tan h k d. So, the speed at which a wave moves in the direction of propagation as a function of what a depth can be given by equation. Earlier we obtained the velocity potential at constant depth but now we have found out the speed of the wave as a function of water depth.

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$$C = \sqrt{\frac{gL}{2\pi} \tanh kd} \quad (2.31)$$

$$\text{Or } L = \frac{gT^2}{2\pi} \tanh kd \quad (2.32)$$

• Since the unknown 'L' occurs on both sides (Implicit Eq.) of Eq. (2.32), it has to be solved by trial and error.

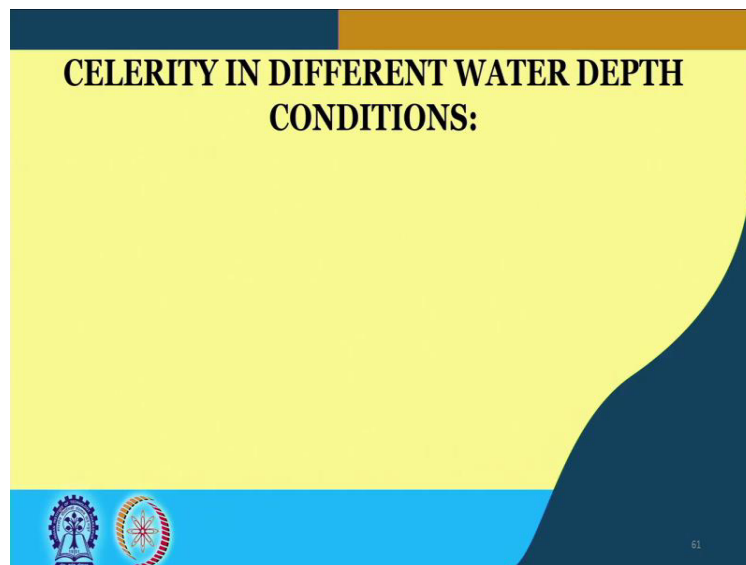
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Since  $C = L$  by  $T$ , we get  $C =$  we just writing it in a different form. So  $C$  can be written as  $gL$  by  $2\pi \tanh kd$ , or  $L = gT^2$  by  $2\pi \tanh kd$ . Since the unknown  $L$  occurs on both side of equation 2.32 it has to be solved by trial and error because  $k$  is nothing but  $2\pi$  by  $L$ . So, if you this, you want to, so, this is the equation dispersion famous dispersion equation, sigma squared =

$g k \tan h k d$  which you must remember and doing some no manipulation here and there we can also write  $L = gT^2 \text{ square by } 2 \pi \tan h k d$ .

And this is also a different form of dispersion relationship and because this  $L$  appears on both sides of this equation 2.32 this equation if needs to be solved will be solved by trial and error method.

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So, I will stop at this point in this lecture and when we start the next lecture, we are going to study the celerity in different water depth conditions. So, thank you so much for listening in this lecture and I will see you in the next class.