


• Using Eq. 14 in Eq. 15, we get $l_m = \kappa y$ — (14)

$$\tau = \rho(\kappa y)^2 \left(\frac{du}{dy} \right)^2$$

or

$$\frac{du}{dy} = \frac{1}{\kappa y} \sqrt{\frac{\tau}{\rho}} \quad (\text{Eq. 16}) \checkmark$$

• For small values of y , it can be assumed that $\tau = \tau_0$ (where τ_0 is the shear stress at the pipe wall and can be assumed to be a constant)



So, if we use equation 14 in equation 15, what was equation 14? l_m was κy , this was what we said in equation number 14. So, we can simply write

$$\tau = \rho(\kappa y)^2 \left(\frac{du}{dy} \right)^2$$

or we can simply write

$$\frac{du}{dy} = \frac{1}{\kappa y} \sqrt{\frac{\tau}{\rho}}$$

and this is equation number 16, very simple. So, for small values of y it can be assumed. So, if the y is very small we can assume that τ is equal to τ_0 , where τ_0 is the shear stress at the pipe wall and can be assumed to be a constant. So, at the wall the shear stress is assumed to be constant and equal to τ_0 .

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• Substituting $\tau = \tau_0$ in Eq. 16, we obtain

$$\frac{du}{dy} = \frac{1}{\kappa y} \sqrt{\frac{\tau_0}{\rho}}$$

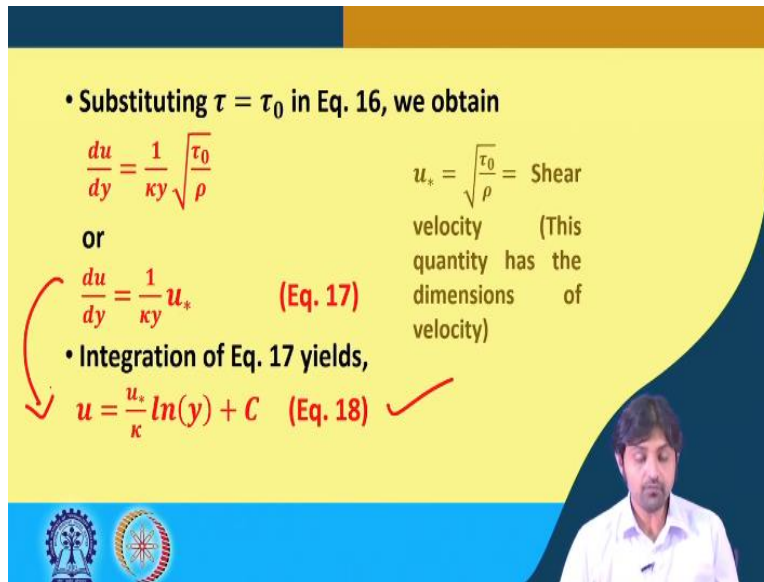
or

$$\frac{du}{dy} = \frac{1}{\kappa y} u_* \quad (\text{Eq. 17})$$

• Integration of Eq. 17 yields,

$$u = \frac{u_*}{\kappa} \ln(y) + C \quad (\text{Eq. 18})$$

$u_* = \sqrt{\frac{\tau_0}{\rho}} = \text{Shear velocity}$ (This quantity has the dimensions of velocity)



And therefore, what we can say, if we substitute tau is equal to tau not in equation 16, we can obtain

$$\frac{du}{dy} = \frac{1}{\kappa y} \sqrt{\frac{\tau_0}{\rho}}$$

or du / dy, this quantity actually tau not can be written as rho. So, but the catch here is, what is the catch? We have considered small value of y. So, du / dy can be written as, 1 / kappa y and under root tau / rho is rho u *. So, it becomes

$$\frac{du}{dy} = \frac{1}{\kappa y} u_*$$

and this u * under root tau not / rho is the sheer velocity and this has the dimension of velocity.

And if you integrate the equation number 17, so, what we can get is, simple integration, it will get

$$u = \frac{u_*}{\kappa} \ln(y) + C$$

. This is very simple integration, from here to here, you can attempt it.

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- Using the boundary condition, $u(y = R) = u_{max}$ (where R is the radius of the pipe) in Eq. 18

$$u = u_{max} + \frac{u_*}{\kappa} [\ln(y) - \ln(R)]$$

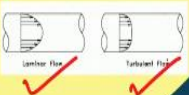

or

$$u = u_{max} + \frac{u_*}{\kappa} \ln\left(\frac{y}{R}\right) \quad (\text{Eq. 19})$$

- Substituting $\kappa = 0.4$ in Eq. 19

$$u = u_{max} + 2.5u_* \ln\left(\frac{y}{R}\right) \quad (\text{Eq. 20})$$

Logarithmic velocity profile

Then using the boundary conditions, what are the boundary conditions? So, u at y is equal to R , where R is the radius of the pipe. We will get, u is equal to u_{max} . That is what we have seen at the center line of the pipe the velocity is going to be the maximum. So, if we use this boundary condition u at y is equal to R is u_{max} , we can get u is equal to, you know, we put u_{max} here, y will be R and therefore, we can obtain C .

C will be $u_{max} - \frac{u_*}{\kappa} \ln R$ and if we substitute this as C , then we can get equation

$$u = u_{max} + \frac{u_*}{\kappa} [\ln(y) - \ln(R)]$$

or

$$u = u_{max} + \frac{u_*}{\kappa} \ln\left(\frac{y}{R}\right)$$

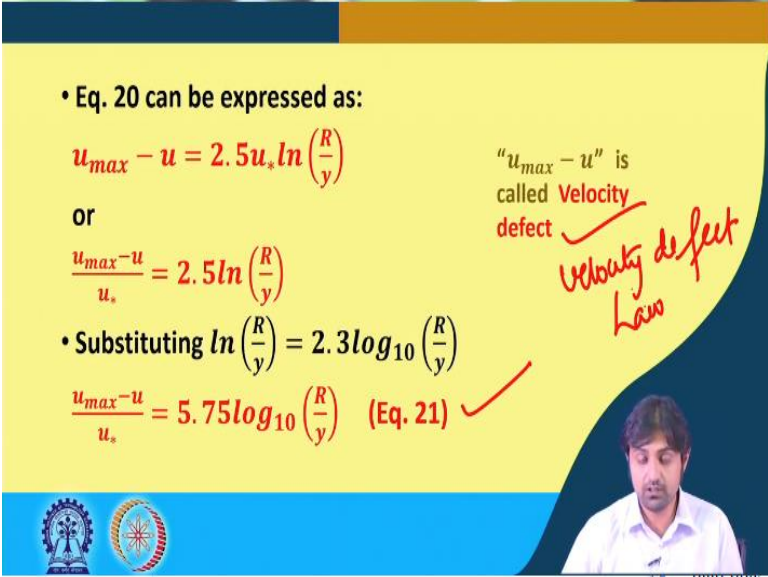
or if when we substitute κ as 0.4, we can get

$$u = u_{max} + 2.5u_* \ln\left(\frac{y}{R}\right)$$

and this is equation number 20. This is just simple manipulation and as you can see we have derived a logarithmic velocity profile starting with the Prandtl mixing length theory for turbulent fluid flow.

So, laminar flow was something like this, a parabolic profile. Here, a profile is little different u is u_{\max} plus a logarithmic profile. So, it looks like something like this. Now, the equation 20 this equation 20 can be expressed as,

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• Eq. 20 can be expressed as:

$$u_{\max} - u = 2.5 u_* \ln \left(\frac{R}{y} \right)$$

or

$$\frac{u_{\max} - u}{u_*} = 2.5 \ln \left(\frac{R}{y} \right)$$

• Substituting $\ln \left(\frac{R}{y} \right) = 2.3 \log_{10} \left(\frac{R}{y} \right)$

$$\frac{u_{\max} - u}{u_*} = 5.75 \log_{10} \left(\frac{R}{y} \right) \quad (\text{Eq. 21})$$

" $u_{\max} - u$ " is called **Velocity defect**

Velocity defect Law

u_{\max} , so, what we do is, we bring u on the other side. So, we bring u on this side and we take this whole side component this side, then what the result is, $u_{\max} - u$ because u_{\max} will always be larger than u is equal to $2.5 u_* \ln R / y$. y will always be less than R or we bring u frictional velocity down, then we get u_{\max} , you bring it down here by dividing then you get $u_{\max} - u / u_*$ equals to $2.5 \ln R / y$. And, so, this is \ln .

So, we can put it in form of \log . This is simple manipulation, we can get $u_{\max} - u / u_*$ is equal to $5.75 \log$ to the base 10 R / y . $u_{\max} - u$ is called the velocity defect or velocity defect law, this is velocity defect law. This is just simple, you know, manipulation of these terms here.

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Problem- 7

- The velocities of water through a pipe of diameter 10 cm are 4 m/s and 3.5 m/s at the center of the pipe and 2 cm from the pipe center, respectively. Considering turbulent flow in the pipe, determine the shear stress at the wall.

(10)

So, now we are going to solve one of the problems, problem number 7. And what it says is, the velocity of water. So, what we have learned in this particular lecture is about the turbulent flow and this problem 7 will help you in solving any problem that is based on this particular concept. So, it says the velocities of water through a pipe of diameter 10 centimeter are 4 meters per second and 3.5 meters per second at the center of the pipe and 2 centimeters from the pipe center, respectively. Considering turbulent flow in pipe, determine the sheer stress at the wall. So, we need to determine tau not. So, let us see how are we going to solve this problem. We are going to have a white screen first.

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Given
 $D = 10 \text{ cm} = 0.1 \text{ m}$
 $R = 0.05 \text{ m}$
 $u_{\text{max}} = 4 \text{ m/s}$ (i.e. at $y=R$)
 $u(r=2 \text{ cm}) = 3.5 \text{ m/s}$
 $y = R - r = 5 - 2 = 3 \text{ cm}$
 $u(y=3) = 3.5 \text{ m/s}$

Now $\frac{u_{\text{max}} - u}{u_*} = 5.75 \log \left(\frac{R}{y} \right)$

$\frac{4 - 3.5}{u_*} = 5.75 \log_{10} \left(\frac{5}{3} \right)$
 $\Rightarrow u_* = 0.392 \text{ m/s}$

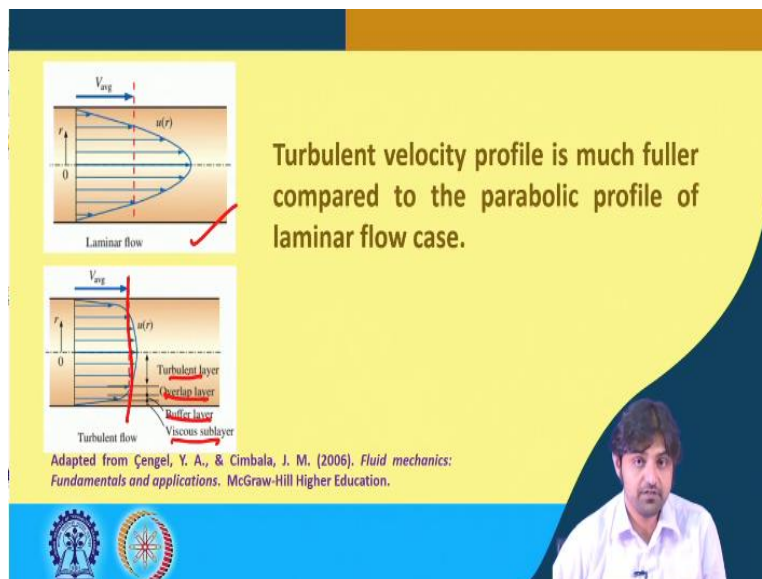
We also know
 $u_* = \sqrt{\frac{\tau_0}{\rho}} \Rightarrow \tau_0 = \rho u_*^2$
 $\tau_0 = 1000 \times (0.392)^2$
 $\tau_0 = 153.6 \text{ N/m}^2$

As always what we do we solve, we write given, diameter is given as 10 centimeter, try to always write down in SI units. So, we write 0.1 meter. So, diameter is 10. So, radius is going to be 0.05 meter. u_{max} is given, is given as 4 meters per second, that is, at y is equal to R . And this is also given, u at r is equal to 2 centimeter is given 3.5 meters per second, that is, y is equal to $R - r$. So, y is going to be $5 - 2$ is equal to 3 centimeter.

So, u at y is equal to 3 is equal to 3.5 meters per second. So, now u_{max} we are using the minus u / u^* was $5.75 \log R / y$. So, substituting the values here, this from here, this equation, $4 - 3.5$ divided by u^* is equal to $5.75 \log_{10} 5 / 3$. This will give us, u^* as 0.392 meters per second. We also know, u^* is under root τ_{not} / ρ or τ_{not} is ρu^* whole square. Therefore, τ_{not} is 1000 and u^* we already got, 0.392 whole square.

So, τ_{not} is coming out to be 153.6 Newton per meter square. This is the solution to the question that we have at hand. So, going back again to the slide, so, what we got was approximately 153 Newton per meter square the sheer stress at the wall.

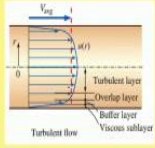
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So, now, the turbulent velocity profile is much fuller compared to the parabolic profile of laminar flow case. Actually this is the flow, this is the true picture, this is a laminar flow that we have seen before. But below is, this is the V average and the velocity fluctuates or deviates from these depending upon the flow condition. So, this is the V average line. There are several other

layers, viscous sublayer, buffer layer, overlap layer and turbulent layer. So, as I told you in the last slide, there are different layers, different layers in turbulent flow.

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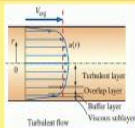


The diagram shows a velocity profile $u(y)$ versus distance from the wall y . The layers are labeled from bottom to top: Viscous sublayer, Buffer layer, Overlap layer, and Turbulent layer. The free stream velocity is U_{∞} and the wall shear stress is τ_w . The turbulent flow is indicated by the dashed line representing the velocity profile.

- Turbulent flow along a wall consist of 4 regions:
 - Viscous sublayer– Thin layer next to the wall where viscous effects are dominant; The velocity profile is almost linear.
 - Buffer layer– Though turbulent effects are becoming significant, the viscous effects are still dominating.

And we are going to talk about that. Turbulent flow along a wall consists of 4 regions. Viscous sublayer, this layer is thin layer next to the wall. So, this is the closest to the wall where the viscous effects are dominant and the velocity profile is all most linear. So, in viscous sub layer the viscous effects are dominant and the velocity profile is linear. In the buffer layer, though turbulent effects are becoming significant, the viscous effects are still dominating.

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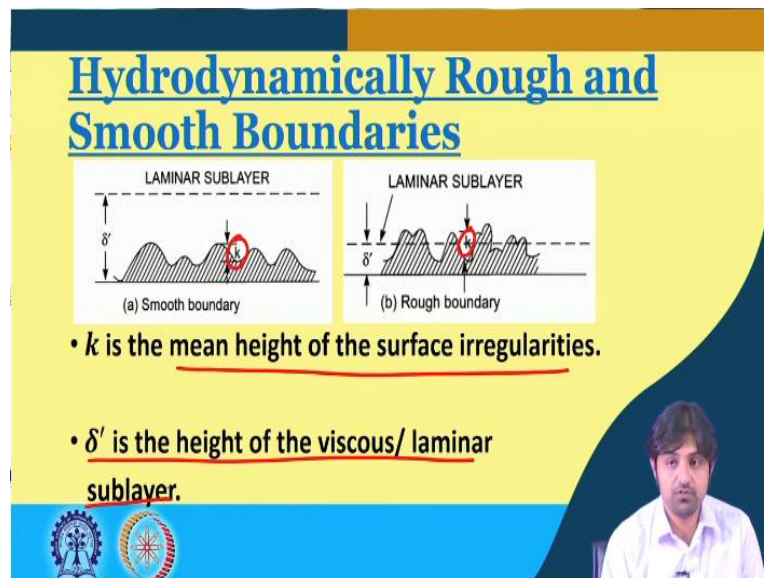


The diagram shows a velocity profile $u(y)$ versus distance from the wall y . The layers are labeled from bottom to top: Viscous sublayer, Buffer layer, Overlap layer, and Turbulent layer. The free stream velocity is U_{∞} and the wall shear stress is τ_w . The turbulent flow is indicated by the dashed line representing the velocity profile.

- Overlap layer– Turbulent effects are much more significant but still not dominant.
- Turbulent layer– Turbulent effects dominate over viscous effects.

In the overlap layer, the turbulent effects are much more significant but still not dominant, in the overlap layer. In the turbulent layer, the turbulent effects dominate over these viscous effects.


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Now, when it comes to these beds and these regimes, some of the important terms that are there is hydro dynamically rough and smooth boundaries. So, this is the, if you see, there is a term called k . Here, if in here, so, k here is the mean height of the surface irregularities. We talked in the beginning that the turbulence could occur due to the presence of irregularities on the surface. So, let us say, the mean height of the surface irregularities is k . And δ' , for example, is the height of viscous or laminar sublayer, the first layer that we talked, the viscous sub layer that was where the velocity profile was almost linear.

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- Outside the laminar sublayer, the flow is turbulent.
- Eddies present in the turbulent zone try to penetrate the laminar sublayer and interact with the boundary.
- When $k \ll \delta'$, the eddies are unable to reach the surface irregularities. **Smooth Boundary**




So, outside the laminar sublayer the flow is turbulent, that is, what we have talked about. Eddies present in the turbulent zone try to penetrate the laminar sublayer and interact with the boundary. But when the surface irregularities are much smaller than δ' , the height of the viscous sublayer, the eddies are unable to reach the surface irregularities when the roughness height is much less. Therefore, we define that boundary as smooth boundary.

So, smooth boundary are the one, where the thickness of the viscous sublayer is much larger than the surface irregularities. We will see, what those surface regularities here, represented by k .

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- When $k \gg \delta'$, the irregularities are above the laminar sublayer leading to the interaction of the eddies with the surface irregularities. **Rough Boundary**
- From Nikuradse's experiments:
 - If $\frac{k}{\delta'} < 0.25$, the boundary is **smooth.**
 - If $\frac{k}{\delta'} > 6$, the boundary is **rough.**



When k is much larger than the delta dash, that is, the thickness of viscous sub layer, the irregularities are above the laminar sublayer leading to the interaction of eddies with the surface irregularities and therefore, these are called rough boundaries. From Nikuradse's roughness, k / δ^+ if it is less than 0.25. So, these values which we are going to talk about, has been derived from experiments by Nikuradse. Nikuradse said if k which is the height of the irregularity is divided by the thickness of viscous sublayer is less than 0.25, the boundary is smooth, if k / δ^+ is greater than 6, the boundary is for sure rough.

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- If $0.25 < \frac{k}{\delta^+} < 6$, the boundary is in transition.
- In terms of roughness Reynolds number $\frac{u_* k}{\nu}$
- If $\frac{u_* k}{\nu} < 4$, the boundary is smooth.
- If $\frac{u_* k}{\nu} > 100$, the boundary is rough.
- If $4 < \frac{u_* k}{\nu} < 100$, the boundary is in transition.

Handwritten note: $Re_k = \frac{u_ k}{\nu}$*

But if it lies in between 0.25 and 6, the boundary is transitional. In terms of roughness Reynolds number, so actually, there is something called roughness Reynolds number that is dependent upon k the height of the irregularities. So, in terms of roughness Reynolds number, if this Reynolds number is less than the 4, the boundaries is smooth, if it is more than the 100 then the boundary is rough and if it is lies between 4 and 100 the boundaries is transitional.

So, either we can calculate it in terms of k / δ^+ , where k is this height of the irregularities and δ^+ is the viscous sublayer or more it is more easy to calculate, $u_* k / \nu$. If this is less than 4, it is smooth, if it is more than 100 then rough otherwise in between it is a transitional boundary.

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Problem- 8

- A pipeline carrying water has average height of irregularities projecting from the surface of the boundary of the pipe as 0.15 mm. What type of boundary is it? The shear stress at the pipe wall is 4.9 N/m^2 and the kinematic viscosity of water is 0.01 stokes.

$$Re_k = 10.5 \\ \Rightarrow \text{Transitional boundary}$$

Now, we will solve one problem about this particular concept. So, the question is, a pipeline carrying water has average height of irregularities projecting from the surface of the boundary of the pipe as 0.15 millimeter. What type of boundary it is? We have to estimate the rough or smooth or transitional boundary. The shear stress at the pipe wall is 4.9 Newton per meter square and the kinematic viscosity is 0.01 Stokes. So, shear stress at the wall is given. So, we will be able to calculate u^* from here. But better that we go and start doing the problems as we have been doing.

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Given: $k = 0.15 \text{ mm} = 0.15 \times 10^{-3} \text{ m}$
 $\tau_0 = 4.9 \text{ N/m}^2$
 $\nu = 0.01 \times 10^{-4} \text{ m}^2/\text{s}$
 $u_* = \sqrt{\frac{\tau_0}{\rho}} = \sqrt{\frac{4.9}{1000}} = 0.07 \text{ m/s}$
Roughness Reynolds number $Re_k = \frac{u_* k}{\nu}$
 $Re_k = \frac{0.07 \times 0.15 \times 10^{-3}}{0.01 \times 10^{-4}} = 10.5$
 $as Re_k < 100 \Rightarrow \text{boundary is transitional}$

So, we have to write the things that we it has been given to us. k is given as, 0.15 millimeter, it is always a good habit to write it into SI unit into 10 to the power - 3 meter, τ_0 is actually

given here, 4.9 Newton per meter square and ν is also given $0.01 \text{ into } 10 \text{ to the power minus } 4$ meter square per second. Therefore, we can simply calculate u^* under root τ_{not} / ρ , as I told you and this will come out to be under root $4.9 / 1000$, so, it will come out to be 0.07 meters per second, very simple.

So, best is to calculate the roughness Reynolds number Re^* and that is given as, $u^* k / \nu$. So, Re^* is, u^* is 0.07, k is $0.15 \text{ into } 10 \text{ to the power } - 3$ and ν is $0.01 \text{ into } 10 \text{ to the power } - 4$ and that comes to be 10.5. So, as Re^* lies between 4 and 400, this implies that the boundary is transitional. So, just going back to that screen, so, what we have got is Re^* is 10.5 implying transitional boundary.

So, this is the place where we will end this lecture of ours today and resume in the next lecture and will talk about turbulent flow in smooth pipes. So, I will see you in the next lecture. Thank you.